



This talk is concerned with the vanishing viscosity limit for hyperbolic system of conservation laws. We consider the following parabolic approximation of the hyperbolic system

$$u_t + A(u)u_x = \varepsilon(B(u)u_x)_x \quad \text{for } t > 0, x \in \mathbb{R}, \quad (1)$$

$$u(0, x) = \bar{u}(x) \quad \text{for } x \in \mathbb{R}, \quad (2)$$

where  $u : [0, +\infty) \times \mathbb{R} \rightarrow \mathbb{R}^n$  and  $A, B$  are  $n \times n$  matrices satisfying the following conditions for some  $\mathcal{U} \subset \mathbb{R}^n$ .

- ① **Strict Hyperbolicity:** The matrix  $A(u)$  is  $C^3$  function and has  $n$  distinct eigenvalues  $\lambda_1(u) < \dots < \lambda_n(u)$  for  $u \in \mathcal{U}$ .
- ② The matrix  $B(u)$  is a  $C^2$  function and positive symmetric definite with  $B(u) \geq c_0 \mathbb{I}_n$  for  $u \in \mathcal{U}$  and for some  $c_0 > 0$ .
- ③

$$A(u)B(u) = B(u)A(u) \text{ for all } u \in \mathcal{U}. \quad (3)$$

In particular  $B(u)$  has  $n$  eigenvalues  $(\mu_i(u))_{1 \leq i \leq n}$ .

### Remark

In the sequel we consider  $l_1, \dots, l_n, r_1, \dots, r_n$  as left and right eigenvectors of  $A(u)$  such that

$$\|r_i(u)\| = 1 \text{ and } l_i(u) \cdot r_j(u) = \begin{cases} 1 & \text{if } i = j, \\ 0 & \text{if } i \neq j. \end{cases} \quad (4)$$

## Remark

If we assume that  $A(u) = Df(u)$  with  $f$  a regular function from  $\mathbb{R}^n$  in  $\mathbb{R}^n$ , it is well-known [Glimm, Bressan et al] that the limit vanishing viscosity system which is conservative:

$$\begin{cases} u_t + \partial_x f(u) = 0, \\ u(0, \cdot) = \bar{u}, \end{cases}$$

has a unique global weak solution provided that  $u_0$  is small in  $TV(\mathbb{R})$ . The class of uniqueness select some shock which satisfy the Lax condition.

Generally it is assumed that the fields are linearly degenerate ( $\nabla \lambda_i(u) \cdot r_i(u) = 0$ ) or genuinely non linear ( $\nabla \lambda_i(u) \cdot r_i(u) \neq 0$ ).

We can note that we are not able to prove the existence of global weak solution when the system is not conservative.

## Some questions on the vanishing viscosity limit?

- Q1 Can we obtain the existence of global strong solution for the parabolic system which are uniformly bounded in  $TV(\mathbb{R})$ ?
- Q2 Can we prove that the sequence  $(u_\varepsilon)_{\varepsilon>0}$  converges strongly to  $u$ ? When the system is conservative we can expect that  $u$  is a global weak solution of the limit vanishing viscosity system.
- Q3 Is it true that the limit  $u$  depends on the viscosity coefficients  $B(u)$ ?

## Some answers to the previous equations

- Bianchini-Bressan [01]: When  $B(u) = Id$  and  $A(u) = Df(u)$  then the sequence of solution  $(u_\varepsilon)_{\varepsilon>0}$  is uniformly bounded in  $L^\infty(\mathbb{R}^+, TV(\mathbb{R}))$ . The sequence converges strongly in  $L^1_{loc,t,x}$  to  $u$  even when the system is not conservative.

Note that  $u^\varepsilon(t, x) = u(t/\varepsilon, x/\varepsilon)$  where  $u$  solves the following problem with fix viscosity but scaled initial data,

$$u_t + A(u)u_x = (B(u)_x)_x \text{ and } u(0, x) = \bar{u}(\varepsilon x). \quad (5)$$

Observe that:

$$TV(\bar{u}(\varepsilon \cdot)) = TV(\bar{u}(\cdot)).$$

### Remark

We are reduced to get  $TV$  norm on  $u$ . To do this, we wish to estimate the  $L^1$  norm of  $\partial_x u$ , it is a priori natural to decompose the vector  $u_x$  in the basis  $(r_i(u))_{1 \leq i \leq n}$ :

$$u_x = \sum_{i=1}^n v_i r_i(u).$$

Differentiating (5), we obtain a system of  $n$  evolution equations:

$$v_{i,t} + (\lambda_i v_i)_x - (\mu_i v_i)_{xx} = \phi_i.$$

By maximum principle we deduce that for  $t \geq \hat{t} > 0$ :

$$\|v_i(t, \cdot)\|_{L^1} \leq \|v_i(\hat{t}, \cdot)\|_{L^1} + \int_{\hat{t}}^t \int_{\mathbb{R}} |\phi_i(s, x)| dx ds.$$

Unfortunately in general if we consider a travelling wave solution  $u(t, x) = U_i(x - \lambda t)$  with  $\lim U(x)_{x \rightarrow -\infty} = U^-$  and  $\lim U(x)_{x \rightarrow +\infty} = U^+$  representing a viscous  $i$  shock, we observe that:

$$\int_{\mathbb{R}} |\phi_i(t, x)| dx \neq 0.$$

### Remark

*This is the reason why it is important to choose a basis  $(\tilde{r}_i(u))_{1 \leq i \leq n}$  in a clever way such that  $\phi_i = 0$  when we consider a viscous travelling wave. In particular we wish to have:*

$$\partial_x U_i(x - \lambda t) = v_i \tilde{r}_i(u, v_i, \lambda).$$

*It is what are doing Bianchini, Bressan by using the center manifold theorem.*

## Remark

If each characteristic field of  $A$  belongs to the Temple class, that is, the Lax curves are straight lines and satisfying

$$Dr_i(u) \cdot r_i(u) = 0 \text{ for all } i = 1, 2, \dots, n, \quad (6)$$

then it can be observed that  $u(t, x) = U(x - \sigma t)$  with  $u_x = a(s)r_i(U(s))$  forms a travelling wave when  $U, a, \sigma_i$  satisfying the following system of ODE,

$$U'(s) = a(s)r_i(U(s)),$$

$$a'(s) = \frac{1}{\mu_i(U(s))} [(\lambda_i(U(s)) - \sigma_i)a(s) - r_i(U(s)) \bullet \mu_i],$$

$$\sigma_i'(s) = 0.$$

We can show in particular in this case that setting  $v_i(t, x) = a(x - \sigma t)$ , we obtain:

$$\partial_t v_i + (\lambda_i(U)v_i)_x - (\mu_i(U)v_i)_{xx} = 0.$$

We set  $u_x^i := l_i(u) \cdot u_x$  and we have

$$u_x = \sum_i u_x^i r_i(u). \quad (7)$$

The directional derivative of a function  $g : \mathbb{R}^n \rightarrow \mathbb{R}^n$  is denoted by  $\zeta \bullet g$  for some  $\zeta \in \mathbb{R}^n$ . More precisely,

$$\zeta \bullet g(u) = \lim_{z \rightarrow 0} \frac{g(u + z\zeta) - g(u)}{z}.$$

$$u_t + \sum_i \lambda_i u_x^i r_i = \sum_i (u_x^i B(u) r_i)_x = \sum_i (\mu_i u_x^i r_i)_x = \sum_i (\mu_i u_x^i)_x r_i + \sum_{i,j} \mu_i u_x^i u_x^j r_j \bullet r_i, \quad (8)$$

Hence, after tedious computation we can write an equation on  $\partial_x u$  which gives:

$$\sum_i (u_{xt}^i + (\lambda_i u_x^i)_x - (\mu_i u_x^i)_{xx}) r_i = \sum_{i,j} p_{ij} u_x^i u_x^j + \sum_{i,j} q_{ij} u_{xx}^i u_x^j + \sum_{i,j,k} s_{ijk} u_x^i u_x^j u_x^k,$$

where  $p_{ij}, q_{ij}$  and  $s_{ijk}$  are defined as follows,

$$\begin{aligned} p_{ij} &= -\lambda_i (r_j \bullet r_i - r_i \bullet r_j), \\ q_{ij} &= 2\mu_i r_j \bullet r_i + (\mu_i - \mu_j) r_j \bullet r_i, \\ s_{ijk} &= 2(r_j \bullet \mu_i) r_k \bullet r_i + \mu_i (r_k \bullet (r_j \bullet r_i) - (r_k \bullet r_i) \bullet r_j) - (r_k \bullet \mu_j) r_j \bullet r_i. \end{aligned}$$

Furthermore, we set  $p_{jk}^i := l_i \cdot p_{jk}$ ,  $q_{jk}^i := l_i \cdot q_{jk}$  and  $s_{jkl}^i := l_i \cdot s_{jkl}$ . Writing  $v_i = u_x^i$ , we get

$$\begin{aligned} v_{i,t} + (\lambda_i v_i)_x - (\mu_i v_i)_{xx} &= \sum_{j,k} p_{jk}^i v_j v_k + \sum_{j,k} q_{jk}^i v_{j,x} v_k + \sum_{j,k,l} s_{jkl}^i v_j v_k v_l \\ &=: \phi_i(u, v_1, \dots, v_n). \end{aligned}$$

We note that  $p_{kk}^i = q_{kk}^i = s_{kkk}^i = 0$  for all  $i, k$  due to the assumption  $r_k \bullet r_k = 0$  for all  $k$ .

### Triangular system

Let us consider now the following triangular system:

$$\begin{cases} u_{1,t} + (f(u_1))_x = 0, \\ u_{2,t} + (g(u_1, u_2))_x = 0. \end{cases}$$

We consider the corresponding viscosity approximation

$$\begin{cases} u_{1,t} + (f(u_1))_x = \alpha_1 u_{1,xx}, \\ u_{2,t} + (g(u_1, u_2))_x = [(\beta(u_1, u_2) u_{1,x})_x + \alpha_2 u_{2,xx}]. \end{cases}$$

We can write (7)–(7) in the following form

$$u_t + A(u)u_x = (B(u)u_x)_x, \tag{9}$$



where  $A, B$  are defined as follows

$$A(u) = \begin{pmatrix} f'(u_1) & 0 \\ \frac{\partial g}{\partial u_1} & \frac{\partial g}{\partial u_2} \end{pmatrix} \text{ and } B(u) = \begin{pmatrix} \alpha_1 & 0 \\ \beta(u) & \alpha_2 \end{pmatrix}. \quad (10)$$

We assume that  $\beta(u)$  satisfies the following condition

$$\beta(u) = (\alpha_1 - \alpha_2) \frac{\frac{\partial g}{\partial u_1}}{f'(u_1) - \frac{\partial g}{\partial u_2}}. \quad (11)$$

which corresponds to  $A(u)B(u) = B(u)A(u)$ . Furthermore we have:

$$r_1(u) = \begin{pmatrix} 1 \\ h(u) \end{pmatrix} \text{ and } r_2(u) = \begin{pmatrix} 0 \\ 1 \end{pmatrix}. \quad (12)$$

## Viscous travelling wave

We would like to decompose  $u_x$  in terms to travelling waves of (9). Let  $u(t, x) = U(x - \sigma_1 t)$  be a travelling wave corresponding to 1-family. Then we have the following ordinary differential system

$$\left. \begin{aligned} \dot{u} &= v, \\ \dot{v} &= B^{-1}(u)(A(u) - \sigma)v - B^{-1}(u)(v \cdot DB(u))v, \\ \dot{\sigma} &= 0. \end{aligned} \right\} \quad (13)$$

We note that  $P_1^* := (u^*, 0, \lambda_1(u^*))$  are equilibrium  $P$  points. We linearize near the point  $P_1^*$  and get

$$\dot{u} = v,$$

We define  $V_i, 1 \leq i \leq 2$  as follows

$$v = \sum_j V_j r_j^*, \quad V_j := l_j^* \cdot v. \quad (15)$$

The center subspace will look like

$$\mathcal{N}_1 := \{(u, v, \sigma) \in \mathbb{R}^n \times \mathbb{R}^n \times \mathbb{R}; V_j = 0, j \neq i\}. \quad (16)$$

Note that  $\dim(\mathcal{N}_1) = 4$ , by Center Manifold Theorem, there exists a smooth manifold  $\mathcal{M}_1 \subset \mathbb{R}^5$  which is tangent to  $\mathcal{N}_1$  at  $P_i^*$ .

Furthermore,  $\mathcal{M}_1$  has dimension 4 and is locally invariant under the flow of (13).

We can write

$$V_2 = \varphi_2(u, V_1, \sigma). \quad (17)$$

We can assume that  $\varphi_2$  is defined on the domain

$$\mathcal{D}_1 := \{|u - u^*| < \varepsilon, |V_1| < \varepsilon, |\sigma - \lambda_1(u^*)| < \varepsilon\}. \quad (18)$$

Note that equilibrium points  $(u, 0, \sigma)$  with  $|u - u^*| < \varepsilon, |\sigma - \lambda_1(u^*)| < \varepsilon$  lie in  $\mathcal{M}_1$  we have

$$\varphi_2(u, 0, \sigma) = 0 \text{ for all } . \quad (19)$$

Hence, we may write

$$\varphi_2(u, V_1, \sigma) = \psi_2(u, V_1, \sigma)V_1, \quad (20)$$

for some  $\psi_2$ .

It implies that:

$$v = V_1(r_1^* + \psi_2(u, V_1, \sigma)r_2^*)$$

We deduce that:

$$v = V_1 \langle l_1(u), (r_1^* + \psi_2(u, V_1, \sigma)r_2^*) \rangle r_1(u) + V_1 \langle l_2(u), (r_1^* + \psi_2(u, V_1, \sigma)r_2^*) \rangle r_2(u).$$

Now, we would like to make a change of coordinates  $V_k \mapsto \tilde{V}_k$  as follows

$$\tilde{V}_k = \langle v, l_k(u) \rangle \quad (21)$$

Therefore, for any point  $(u, v, \sigma) \in \mathcal{M}_1$  we can write

$$v = \tilde{V}_1 \left( r_1(u) + \tilde{\psi}_2(u, \tilde{V}_1, \sigma)r_2(u) \right) =: \tilde{V}_1 \tilde{r}_1(u, \tilde{V}_1, \sigma). \quad (22)$$

We note that the gradient of 1-family travelling waves can be written under the following form

$$u_x = v_1 \tilde{r}_1 \text{ where } \tilde{r}_1 = \begin{pmatrix} 1 \\ s(u, v_1, \sigma_1) \end{pmatrix} \quad (23)$$

with

$$s(u, v_1, \sigma_1) = \tilde{\psi}_2(u, v_1, \sigma_1) + \frac{\frac{\partial g}{\partial u_1}}{f'(u_1) - \frac{\partial g}{\partial u_2}}.$$

Furthermore we can check that:

$$\tilde{r}_{1,\sigma} = \mathcal{O}(1)v_1, \quad \tilde{r}_{1,\sigma\sigma} = \mathcal{O}(1)v_1, \quad \tilde{r}_1 \bullet \tilde{r}_{1,\sigma} = \mathcal{O}(1)v_1. \quad (24)$$

We want to have a decomposition of  $u_x$  as

$$u_x = v_1 \tilde{r}_1(u, v_1, \sigma_1) + v_2 r_2(u). \quad (25)$$

We observe that  $B\tilde{r}_1 = \alpha_1 \tilde{r}_1 + s_1 r_2$  with  $s_1(u, v_1, \sigma_1)$  a function depending on  $(u, v_1, \sigma_1)$ . Set  $w_1 = \alpha_1 v_{1,x} - \lambda_1 v_1$  the effective flux, after tedious computation we obtain that:

$$\begin{aligned} (B(u)u_x)_x - A(u)u_x &= w_1 [\tilde{r}_1 + v_1 \tilde{r}_{1,v}] + (\alpha_2 v_{2,x} - \lambda_2 v_2) r_2 \\ &+ \frac{1}{\alpha_1} (w_1 + \sigma_1 v_1) (s_1 + v_1 s_{1,v}) r_2 + v_1 \sigma_{1,x} s_{1,\sigma} r_2 \\ &+ v_1^2 \sigma_1 \tilde{r}_{1,v} + \alpha_1 \sigma_{1,x} v_1 \tilde{r}_{1,\sigma} + v_1 v_2 r_2 \bullet s_1 r_2 + \alpha_1 v_1 v_2 r_2 \bullet \tilde{r}_1. \end{aligned}$$

We need now to specify the choice of  $\sigma_1$ . Assume that we consider a travelling wave such that  $u_x = v_1 \tilde{r}_1$  then from the previous equation we should have:

$$\begin{aligned} u_t &= w_1 \tilde{r}_1 + \frac{1}{\alpha_1} (w_1 + \sigma_1 v_1) (s_1 + v_1 s_{1,v}) r_2 \\ &+ (w_1 + v_1 \sigma_1) v_1 \tilde{r}_{1,v}. \end{aligned}$$

Since we have  $u_t = -\sigma_1 u_x$ , we must choose  $\sigma_1$  such that  $\sigma_1 = -\frac{w_1}{v_1}$ .

Since  $\sigma_1$  must live in a neighborhood of  $\lambda_1(u^*)$ , we set

$$\sigma_1 = \lambda_1(u^*) + \theta\left(-\frac{w_1}{v_1} - \lambda_1(u^*)\right),$$

with

$$\theta(s) = \begin{cases} s & \text{if } |s| \leq \frac{\delta_1}{2}, \\ 0 & \text{if } |s| \geq \delta_1, \end{cases} \quad |\theta'| \leq 1 \text{ and } |\theta''| \leq 4/\delta_1. \quad (26)$$

Let us assume now that we consider a solution  $u_x$  of our system such that:

$$u_x = v_1 \tilde{r}_1(u, v_1, \sigma_1) + v_2 r_2(u). \quad (27)$$

If we consider the first coordinate of  $u_x$  and due to the form of  $\tilde{r}_1$ , we observe that  $v_1 = u_{1,x}$ . After computations, we can show now that  $(v_1, v_2)$

$$\begin{cases} v_{1,t} + (\lambda_1 v_1)_x - \alpha_1 v_{1,xx} = 0, \\ v_{2,t} + (\lambda_2 v_2)_x - \alpha_2 v_{2,xx} = \phi_2, \end{cases} \quad (28)$$

with:

$$\begin{aligned} \phi_2 = & O(1)|v_{1,x}(w_1 + \sigma_1 v_1)| \quad \text{wrong speed} \\ & + O(1)|w_{1,x}v_1 - v_{1,x}w_1| \quad \text{change in strength} \\ & + O(1)|v_1[v_1\left(\frac{w_1}{v_1}\right)_x]| \chi_{\{x, |\frac{w_1}{v_1}| \leq 3\delta_1\}} \quad \text{change in speed} \\ & + O(1)[|v_1 v_2| + |v_{1,x} v_2|] \quad \text{transversal interactions} \\ & + O(1)\left[ \left| \frac{1}{\alpha_1} ((w_1 + \sigma_1 v_1)(s_1 + v_1 s_{1,v}))_x \right| + \left| (v_1 \sigma_{1,x} s_{1,\sigma})_x \right| \right] \end{aligned}$$

## Theorem

Consider the Cauchy problem hyperbolic system with viscosity,

$$u_t + A(u)u_x = \varepsilon(B(u)u_x)_x, \quad u(0, x) = \bar{u}(x). \quad (29)$$

There exists  $L_1, L_2, L_3 > 0$  and  $\delta_0 > 0$  such that the following holds. If  $\bar{u}$  satisfies

$$TV(\bar{u}) \leq \delta_0 \text{ and } \lim_{x \rightarrow \mathbb{R}^-} \bar{u}(x) \in K, \quad (30)$$

for some compact set  $K \subset \mathcal{U}$  then there exists unique solution  $u^\varepsilon$  to the Cauchy problem (29) and it satisfies the following properties

$$TV(u^\varepsilon(t)) \leq L_1 TV(\bar{u}), \quad (31)$$

$$\|u^\varepsilon(t) - v^\varepsilon(t)\|_{L^1} \leq L_2 \|\bar{u} - \bar{v}\|_{L^1}, \quad (32)$$

$$\|u^\varepsilon(t) - u^\varepsilon(s)\|_{L^1} \leq L_3 \left( |t - s| + \sqrt{\varepsilon} |\sqrt{t} - \sqrt{s}| \right), \quad (33)$$

where  $v^\varepsilon$  is the unique solution corresponding to  $\bar{v}$  satisfying (30).

Furthermore, when  $A = Df$  for some  $f \in C^1$ , as  $\varepsilon \rightarrow 0$  (up to a subsequence),  $u^\varepsilon \rightarrow u$  in  $L^1_{loc}$  with  $u$  a solution to hyperbolic system (5).

## Theorem

Then for every compact set  $K \subset \mathcal{U}$  there exist  $L_1, L_2, \delta_0$ , a closed domain  $\mathcal{D} \subset L^1_{loc}(\mathbb{R})$  and a semigroup  $S : [0, \infty) \times \mathcal{D} \rightarrow \mathcal{D}$  satisfying the following properties.

- ① Every function  $\bar{u}$  verifying (30) belongs to  $\mathcal{D}$ .
- ② For any  $\bar{u}, \bar{v} \in \mathcal{D}$  with  $\bar{u} - \bar{v} \in L^1$ ,

$$\|S_{t_1}(\bar{u}) - S_{t_2}(\bar{v})\|_{L^1} \leq L\|\bar{u} - \bar{v}\|_{L^1} + L'|t_1 - t_2| \text{ for any } t_1, t_2 \geq 0, \quad (34)$$

for some constants  $L, L'$  which are depending only on  $\mathcal{D}$ .

- ③ For any piece-wise constant initial data  $\bar{u} \in \mathcal{D}$  there exists  $\tau > 0$  such that the following holds. For  $t \in [0, \tau]$ ,  $S_t$  coincides with the solution constructed by gluing the Riemann problem solutions arising at each jump point.
- ④ For each  $\bar{u} \in \mathcal{D}$ ,  $t \mapsto S_t(\bar{u})$  is the unique limit of the sequence  $u^{\varepsilon_k}(t, \cdot)$  in  $L^1_{loc}$  for any  $\varepsilon_k \rightarrow 0$  where  $u^{\varepsilon_k}(t, \cdot)$  solves (1) with initial data  $\bar{u}$ .

## Remark

Let  $S_t^I$  be the semigroup constructed by Bianchini-Bressan for (5). Due to the characterization S3, we conclude that the semigroup  $S_t^B$  constructed as above coincides with  $S_t^I$ .

## Parabolic estimates

### Proposition

Let  $u$  be a solution to the equation (5) satisfying

$$\|u_x(t, \cdot)\|_{L^1} \leq \delta_0 \text{ for all } t \in [0, \hat{t}] \text{ where } \hat{t} := \left(\frac{1}{C\delta_0}\right)^2, \quad (35)$$

for some  $\delta_0 < 1$  and  $C > 0$ . Then we have

$$\begin{aligned} \|u_{xx}(t, \cdot)\|_{L^1} &\leq \frac{2\kappa\kappa_1^2\kappa_P^2\delta_0}{\sqrt{t}} \\ \|u_{xxx}(t, \cdot)\|_{L^1} &\leq \frac{5\kappa\kappa_1^2\kappa_P^2\delta_0}{\sqrt{t}} \\ \|u_{xxx}(t, \cdot)\|_{L^\infty} &\leq \frac{16\kappa\kappa_1^2\kappa_P^2\delta_0}{\sqrt{t}} \end{aligned} \quad (36)$$



First we have:

$$(u_x)_t + A(u)u_{xx} = B(u)u_{xxx} + (u_x \bullet B(u)u_x)_x - u_x \bullet A(u)u_x. \quad (37)$$

We would like to diagonalize the system, making a change of variable  $v = P(u)u_x$  we get:

$$\begin{aligned} v_t + A_1(u)v_x &= B_1(u)v_{xx} - B_1((P^{-1}(u)v) \bullet PP^{-1}v)_x - B_1[(P^{-1}v) \bullet P(P^{-1}v)_x] \\ &\quad + u_t \bullet P(u)P^{-1}v + A_1(u)(P^{-1}(u)v \bullet P(u)P^{-1}v) \\ &\quad + P(P^{-1}v \bullet B(u)P^{-1}v)_x - P(P^{-1}v) \bullet A(u)P^{-1}v. \end{aligned}$$

with  $A_1 = PAP^{-1}$  and  $B_1 = \text{diag}(\mu_1(u), \dots, \mu_n(u)) = PBP^{-1}$ . Next, we do a change of variable  $v \mapsto \tilde{v}$  such that  $v(t, x) = (\tilde{v}_i(t, X_i(t, x)))$  where  $(X_i)_x = \frac{1}{\sqrt{\mu_i(u)}}$ . Then we have

$$\begin{aligned} \tilde{v}_t + A_2^* \tilde{v}_x &= \tilde{v}_{xx} + \mathcal{T}([A_2^* - A_1^* B_1^{-1/2}(u)] \tilde{v}_x + B_1(P^{-1}v) \bullet B_1^{-1/2} \tilde{v}_x) \\ &\quad - \mathcal{T}(\tilde{v}_{i,x} X_{i,t}) + \mathcal{T}(\mathcal{R}), \end{aligned} \quad (38)$$

where  $A_2^* = A_1^* B_1^{-1/2}(u^*)$  with:

$$\mathcal{T}(f)_i(x) = f_i(X_i(x)) \text{ where } X_i(x) = \int_0^x \frac{1}{\sqrt{\mu_i(u(z))}} dz. \quad (39)$$

We consider  $G$  as the fundamental solution of the following parabolic equation

$$w_t + A_2^* w_x = w_{xx}, \quad (40)$$

where  $A_2^* = A_1^* B_1^{-1/2}(u^*)$ . The function  $G$  satisfies the following estimates

$$\|G(t, \cdot)\|_{L^1} \leq \kappa, \quad \|G_x(t, \cdot)\|_{L^1} \leq \frac{\kappa}{\sqrt{t}}, \quad \|G(t, \cdot)\|_{L^1} \leq \frac{\kappa}{t}, \quad (41)$$

We argue by contradiction. To this end, first we assume that the conclusion does not hold. Due to the assumption of smoothness of initial data, solution is smooth up to a small time and due to the continuity we can assume that there exists a time  $t^*$  such that (47) holds for  $t \in [0, t^*]$  and equality attains at  $t = t^*$ . We can write for  $t \in [0, t^*]$ :

$$\begin{aligned} \tilde{v}_x = & G_x(t/2) \star \tilde{v}(t/2) + \int_{t/2}^t G_x(t-s) \star \left\{ \mathcal{T}([A_2^* - A_1^* B_1^{-1/2}(u)] \tilde{v}_x \right. \\ & \left. + B_1(P^{-1}v) \bullet B_1^{-1/2} \tilde{v}_x) - \mathcal{T}(\tilde{v}_{i,x} X_{i,t}) + \mathcal{T}(\mathcal{R}) \right\} ds. \end{aligned}$$

Furthermore, we observe that

$$\|\tilde{v}\|_{L^1} \leq \kappa_P \kappa_1 \|u_x\|_{L^1}, \quad (42)$$

$$\|\tilde{v}_x\|_{L^1} \leq \kappa_P \kappa_1 \|u_{xx}\|_{L^1}. \quad (43)$$

We get:

$$\begin{aligned} \|\tilde{v}_x(t)\|_{L^1} &\leq \|G_x(t/2)\|_{L^1} \|\tilde{v}(t/2)\|_{L^1} + \int_{t/2}^t \|G_x(t-s)\|_{L^1} \|\mathcal{T}([A_2^* - A_1^* B_1^{-1/2}(u)]\tilde{v}_x)\|_{L^1} ds \\ &\quad + \int_{t/2}^t \|G_x(t-s)\|_{L^1} \|\mathcal{T}(B_1(P^{-1}v) \bullet B_1^{-1/2}\tilde{v}_x) - \mathcal{T}(\tilde{v}_{i,x}X_{i,t}) + \mathcal{T}(\mathcal{R})\|_{L^1} ds \\ &\leq \frac{2\kappa\delta_0}{\sqrt{t}} + \int_{t/2}^t \frac{\kappa\kappa_1}{\sqrt{t-s}} \| [A_2^* - A_1^* B_1^{-1/2}(u)]\tilde{v}_x \|_{L^1} ds \\ &\quad + \int_{t/2}^t \frac{\kappa\kappa_1}{\sqrt{t-s}} \| B_1(P^{-1}v) \bullet B_1^{-1/2}\tilde{v}_x - (\tilde{v}_{i,x}X_{i,t}) + \mathcal{R} \|_{L^1} ds. \end{aligned}$$

We note that

$$\begin{aligned} \|[A_2^* - A_1^* B_1^{-1/2}(u)]\tilde{v}_x\|_{L^1} &\leq \kappa_A \kappa_B \|u - u^*\|_{L^\infty} \|\tilde{v}_x\|_{L^1} \\ &\leq \kappa_A \kappa_B \kappa_P \kappa_1 \|u - u^*\|_{L^\infty} \|u_{xx}\|_{L^1}, \end{aligned}$$

Therefore, we get

$$\begin{aligned} \|\tilde{v}_x(t)\|_{L^1} &\leq \frac{\sqrt{2}\kappa\kappa_1\kappa_P\delta_0}{\sqrt{t}} + 600\kappa_1^6\kappa^3\kappa_A\kappa_B^7\kappa_P^{12} \int_{t/2}^t \frac{1}{\sqrt{t-s}} \left[ \frac{\delta_0^2}{s} + \frac{\delta_0^2}{\sqrt{s}} \right] ds \\ &< \frac{2\kappa\kappa_1\kappa_P\delta_0}{\sqrt{t}}, \end{aligned}$$

which implies

$$\|u_{xx}(t^*, \cdot)\|_{L^1} < \frac{2\kappa\kappa_1^2\kappa_P^2\delta_0}{\sqrt{t^*}}. \quad (44)$$

This contradicts the assumption that equality holds in (47) at  $t = t^*$ .

## Proposition

Let  $\bar{u}$  satisfying

$$TV(\bar{u}) \leq \frac{\delta_0}{4\kappa}.$$

Then  $u$  are well-defined on  $[0, \hat{t}]$  where  $\hat{t}$  is defined as in (46). Moreover, we have

$$\|u_x(t)\|_{L^1} \leq \frac{\delta_0}{2} \text{ for } t \in [0, \hat{t}]. \quad (45)$$

Suppose that there exists a time  $\tau < \hat{t}$  such that  $\|u_x(\tau)\|_{L^1} = \frac{\delta_0}{2}$  and  $\|u_x(t)\|_{L^1} < \frac{\delta_0}{2}$  for all  $t \in [0, \tau]$ . We can write

$$u_t + A(u^*)u_x = B(u^*)u_{xx} + (B(u) - B(u^*))u_{xx} + (A(u^*) - A(u))u_x + u_x \bullet Bu_x.$$

Therefore,

$$\|u_x(\tau)\|_{L^1} \leq \kappa \|u_{0,x}\|_{L^1} + \int_0^\tau \frac{2\kappa\kappa_B}{\sqrt{\tau-s}} \|u_x\|_{L^1} \|u_{xx}\|_{L^1} ds + \int_0^\tau \frac{\kappa\kappa_A}{\sqrt{\tau-s}} \|u_x\|_{L^1}^2 ds.$$

By Proposition 2.1 we get

$$\|u_{xx}(t, \cdot)\|_{L^1} \leq \frac{2\kappa\kappa_1^2\kappa_P^2\delta_0}{\sqrt{t}} \text{ for } t \in [0, \tau].$$

Hence, by the choice of  $\delta_0$

$$\begin{aligned} \|u_x(\tau)\|_{L^1} &\leq \frac{\delta_0}{4} + 2 \int_0^\tau \frac{\kappa}{\sqrt{\tau-s}} \frac{2\kappa\kappa_1^2\kappa_P^2\delta_0}{\sqrt{s}} \frac{\kappa_B\delta_0}{2} ds + \int_0^\tau \frac{\kappa}{\sqrt{\tau-s}} \frac{\kappa_A\delta_0^2}{4} ds \\ &< \frac{\delta_0}{2}. \end{aligned}$$

### Proposition

Let  $T > \hat{t}$  and  $u$  be a solution to the equation (5) satisfying

$$\|u_x(t, \cdot)\|_{L^1} \leq \delta_0 \text{ for all } t \in [0, T] \quad (46)$$

for some  $\delta_0 < 1$  and  $C > 0$ . Then we have for  $t \in [\hat{t}, T]$

$$\begin{aligned} \|u_{xx}(t, \cdot)\|_{L^1} &= O(1)\delta_0^2 \\ \|u_{xxx}(t, \cdot)\|_{L^1} &= O(1)\delta_0^3 \\ \|u_{xxx}(t, \cdot)\|_{L^\infty} &= O(1)\delta_0^4. \end{aligned} \quad (47)$$

## Interaction estimates

## Lemma

Let  $z, z^\#$  be solutions of the two independent scalar equations,

$$z_t + (\lambda(t, x)z)_x - (\mu z)_{xx} = \varphi(t, x), \quad (48)$$

$$z_t^\# + (\lambda^\#(t, x)z^\#)_x - (\mu^\# z^\#)_{xx} = \varphi^\#(t, x), \quad (49)$$

which is valid for  $t \in [0, T]$ . We assume that  $\inf_{t,x} \lambda^\#(t, x) - \sup_{t,x} \lambda(t, x) \geq c > 0$

and  $\|(\mu, \mu^\#)\|_{L^\infty} < \infty$ ,  $\mu, \mu^\# \geq c_0 > 0$ . Then we have

$$\int_0^T \int_{\mathbb{R}} |z(t, x)| |z^\#(t, x)| dx dt \leq \frac{1}{c} E_1 E_2, \quad (50)$$

$$E_1 := \int_{\mathbb{R}} |z(0, x)| dx + \int_0^T \int_{\mathbb{R}} |\varphi(t, x)| dx dt, \quad (51)$$

$$E_2 := \int_{\mathbb{R}} |z^\#(0, x)| dx + \int_0^T \int_{\mathbb{R}} |\varphi^\#(t, x)| dx dt. \quad (52)$$

Set  $c_1 := \|\mu, \mu^\#\|_{L^\infty}$ . Let  $z, z^\#$  be the solution to (48), (49) with  $\varphi = \varphi^\# = 0$ . Consider

$$Q(z, z^\#) := \int \int K(x-y) |z(x)| |z^\#(y)| dx dy, \quad (53)$$

where  $K$  is defined as follows

$$K(s) := \begin{cases} 1/c & \text{if } s \geq 0, \\ 1/ce^{\frac{cs}{2c_1}} & \text{if } s < 0. \end{cases} \quad (54)$$

Now, we can calculate using the fact that  $cK' - 2c_1K''$  is precisely the Dirac masses

$$\begin{aligned} \frac{d}{dt} Q(z(t), z^\#(t)) &= \int \int K(x-y) [sgn(z(x)) z_t(x) |z^\#(y)| + sgn(z^\#(y)) z_t^\#(y) |z(x)|] dx dy \\ &= \int \int K(x-y) \left[ sgn(z(x)) ((\mu(x)z(x))_{xx} - (\lambda z(x))_x) |z^\#(y)| \right. \\ &\quad \left. + sgn(z^\#(y)) ((\mu^\#(y)z^\#(y))_{yy} - (\lambda^\# z^\#(y))_y) |z(x)| \right] dx dy \\ &\leq - \int \int (cK'(x-y) - 2c_1K''(x-y)) |z(x)| |z^\#(y)| dx dy \\ &\leq - \int |z(x)| |z^\#(x)| dx. \end{aligned}$$

$$\int_0^T \int_{\mathbb{R}} |z(t, x)| |z^\#(t, x)| dx dt \leq Q(z(0), z^\#(0)) \leq \frac{1}{c} \|z(0)\|_{L^1} \|z^\#(0)\|_{L^1}. \quad (55)$$



Now, we consider  $z, z^\#$  as solutions of (48), (49) respectively when  $\varphi$  and  $\varphi^\#$  may not be identically 0. Using the representation of the solution in terms of  $\Gamma, \Gamma^\#$  be the fundamental solutions corresponding to the homogeneous system of (48)–(49) we can conclude in a similar way. Indeed we can write

$$z(t, x) = \int_{\mathbb{R}} \Gamma(t, x, 0, y) z(0, y) dy + \int_0^t \int_{\mathbb{R}} \Gamma(t, x, s, y) \varphi(s, y) dy ds. \quad (56)$$

### Lemma

Let  $z, z^\#$  be solutions of (48), (49) respectively and we assume that

$$\int_0^T \int_{\mathbb{R}} |\varphi(t, x)| dx dt \leq \delta_0, \quad \int_0^T \int_{\mathbb{R}} |\varphi^\#(t, x)| dx dt \leq \delta_0, \quad (57)$$

$$\|z(t)\|_{L^1}, \|z^\#(t)\|_{L^1} \leq \delta_0, \quad \|z_x(t)\|_{L^1}, \|z^\#(t)\|_{L^\infty} \leq C_* \delta_0^2, \quad (58)$$

$$\|\lambda_x(t)\|_{L^\infty}, \|\lambda_x(t)\|_{L^1} \leq C_* \delta, \quad \lim_{x \rightarrow -\infty} \lambda(t, x) = 0, \quad (59)$$

for all  $t \in [0, T]$ . Then we have

$$\int_0^T \int_{\mathbb{R}} |z_x(t, x)| |z^\#(t, x)| dx dt = \mathcal{O}(1) \delta_0^2. \quad (60)$$

## Lemma

Let  $z, z^\#$  be solutions of (48), (49) respectively and we assume that

$$\int_0^T \int_{\mathbb{R}} |\varphi(t, x)| dx dt \leq \delta_0, \quad \int_0^T \int_{\mathbb{R}} |\varphi^\#(t, x)| dx dt \leq \delta_0, \quad (61)$$

$$\|z(t)\|_{L^1}, \|z^\#(t)\|_{L^1} \leq \delta_0, \quad \|z_x(t)\|_{L^1}, \|z^\#(t)\|_{L^\infty} \leq C_* \delta_0^2, \quad (62)$$

$$\|\lambda_x(t)\|_{L^\infty}, \|\lambda_x(t)\|_{L^1} \leq C_* \delta, \quad \lim_{x \rightarrow -\infty} \lambda(t, x) = 0, \quad (63)$$

for all  $t \in [0, T]$ . Then we have

$$\int_0^T \int_{\mathbb{R}} |z_x(t, x)| |z^\#(t, x)| dx dt = \mathcal{O}(1) \delta_0^2. \quad (64)$$

Let  $v, w$  be two scalar functions satisfying:

$$\begin{cases} v_t + (\lambda(t, x)v)_x - (\mu v)_{xx} & = \varphi(t, x), \\ w_t + (\lambda(t, x)w)_x - (\mu w)_{xx} & = \varphi^\#(t, x), \end{cases}$$

Considering the functionals as:

$$\begin{aligned} \mathcal{A}(t) &= \frac{1}{2} \int \int_{x < y} |v(t, x)w(t, y) - v(t, y)w(t, x)| dx dy, \\ \mathcal{L}(t) &= \int \sqrt{v^2(t, x) + w^2(t, x)} dx \end{aligned}$$

we have then the following Lemmas.

### Lemma

*The previous functionals satisfies:*

$$\begin{aligned} \frac{d}{dt} \mathcal{A}(t) + \int |v_x(t, x)w(t, x) - w_x(t, x)v(t, x)| dx \\ \|v(t)\|_{L^1} \|\varphi^\#(t)\|_{L^1} + \|w(t)\|_{L^1} \|\varphi(t)\|_{L^1} \end{aligned} \quad (65)$$

$$\frac{d}{dt} \mathcal{L}(t) \leq -C(\delta_1) \int_{|\frac{w}{v}| \leq 3\delta_1} [v(t) |(\frac{w(t)}{v(t)})_x|^2 dx + \|\varphi^\#(t)\|_{L^1} + \|\varphi(t)\|_{L^1}.$$

## TV Bounds

Let us consider an initial data satisfying  $TV(\bar{u}) \leq \frac{\delta_0}{8\sqrt{n}}$  and  $\lim_{x \rightarrow -\infty} u(x) = u^* \in K$ . Then by applying Proposition 2.2, we obtain

$$\|u_x(\hat{t})\|_{L^1(\mathbb{R})} \leq \frac{\delta_0}{4\sqrt{n}}, \quad (66)$$

where  $\hat{t}$  is defined as in (46). To get the total variation bound in  $(\hat{t}, \infty)$  we argue by contradiction as in Bianchini-Bressan. Let  $T$  be defined as follows

$$T := \sup \left\{ \tau; \sum_i \int_{\hat{t}}^{\tau} \int_{\mathbb{R}} |\phi_i(t, x)| dx dt \leq \frac{\delta_0}{2} \right\}. \quad (67)$$

If  $T < +\infty$ , we get a contradiction as follows. From (67), we have

$$\|u_x(t)\|_{L^1} \leq 2\sqrt{n}\|u_x(\hat{t})\|_{L^1} + \frac{\delta_0}{2} \leq \delta_0 \text{ for all } t \in [\hat{t}, T]. \quad (68)$$

By applying 6, we get

$$\int_{\hat{t}}^{\tau} \int_{\mathbb{R}} \left| \sum_{j,k} q_{jk}^i v_{j,x} v_k \right| dx dt = \mathcal{O}(1)\delta_0^2 < \frac{\delta_0}{2} \quad (69)$$

for sufficiently small  $\delta_0 > 0$ .

Indeed we observe that  $\|v_{j,x}(t)\|_{L^1}$ ,  $\|v_i^\#(t)\|_{L^\infty} \leq C_* \delta_0^2$  are satisfied using the proposition 2.3.

### How to deal with the new terms in the case of the triangular system?

We recall that in  $\phi_2$  we have new terms of the form:

$$(v_1 \sigma_{1,x} s_{1,\sigma})_x.$$

Since  $s_{1,\sigma} = O(1)v_1$  and  $\sigma_1 = \lambda_1(u^*) + \theta(-\frac{w_1}{v_1} - \lambda_1(u^*))$ , we have to deal with a term of the form:

$$w_{1,xx} v_1 - v_{1,xx} w_1.$$

First of all we can observe that:

$$w_{1,t} + (\lambda_1(u)w_1)_x - \alpha_1 w_{1,xx} = 0. \quad (70)$$

Furthermore we have  $\lambda_1(u) = f'(u_1)$ , we introduce now a new variable

$$z_1 = \alpha_1 w_{1,x} - \lambda_1(u)w_1.$$

We can check that this new unknown satisfies:

$$z_{1,t} = \alpha_1 z_{1,xx} - (\lambda_1 z_1)_x + \lambda_1'(w_{1,x} v_1 - w_1 v_{1,x}).$$

We observe now that:

$$z_{1,x}v_1 - v_{1,x}z_1 = w_{1,xx}v_1 - w_1v_{1,xx} + 2\lambda_1(w_1v_{1,x} - w_{1,x}v_1). \quad (71)$$

Using now (65), we deduce that:

$$\begin{aligned} & \int_{\hat{t}}^{\tau} \int_{\mathbb{R}} |w_{1,xx}v_1 - w_1v_{1,xx}| \, dxdt \\ & \leq \int_{\hat{t}}^{\tau} \int_{\mathbb{R}} |z_{1,x}v_1 - v_{1,x}z_1| \, dxdt + 2\|\lambda_1\|_{L^\infty} \int_{\hat{t}}^{\tau} \int_{\mathbb{R}} |w_{1,x}v_1 - v_{1,x}w_1| \, dxdt \\ & \leq \|z_1(\hat{t})\|_{L^1} \|v_1(\hat{t})\|_{L^1} + \|v_1\|_{L^\infty((\hat{t},\tau);L^\infty(\mathbb{R}))} \|\lambda_1'(w_{1,x}v_1 - w_1v_{1,x})\|_{L^1([\hat{t},\tau] \times \mathbb{R})} \\ & \quad + 2\|\lambda_1\|_{L^\infty} \|w_{1,x}v_1 - v_{1,x}w_1\|_{L^1([\hat{t},\tau] \times \mathbb{R})}. \end{aligned}$$

Using again (65), we conclude since:

$$\|w_{1,x}v_1 - v_{1,x}w_1\|_{L^1([\hat{t},\tau] \times \mathbb{R})} = 0(1)\delta_0^2.$$

Hence,  $T$  is not the supremum defined as in (67). Hence,  $\int_{\hat{t}}^{\tau} \int_{\mathbb{R}} |\phi_i(t,x)| \, dxdt \leq \frac{\delta_0}{2}$

for all  $t > \hat{t}$ . Subsequently, we obtain for all  $t \geq 0$ :

$$\|u_x(t)\|_{L^1} \leq \delta_0.$$

## Stability Estimates

Let  $\bar{u}^\theta$  be the initial defined as follows

$$\bar{u}^\theta := \theta\bar{u} + (1-\theta)\bar{v} \text{ for some } \bar{u}, \bar{v} \in \mathcal{D}. \quad (72)$$

Let  $u^\theta$  be the solution associated to initial data  $\bar{u}^\theta$ . Then taking derivative with

We can now write after tedious computations

$$\begin{aligned} & \sum_i (h_{i,t} + (\lambda_i h_i)_x - (\mu_i h_i)_{xx}) r_i \\ &= \sum_{i,j} \hat{p}_{ij} h_i v_j + \sum_{i,j,k} \hat{q}_{ijk} h_i v_j v_k + \sum_{i,j} \hat{s}_{ij} h_{i,x} v_j + \sum_{i,j} \hat{w}_{ij} h_i v_{j,x}, \end{aligned}$$

where  $\hat{p}_{ij}, \hat{q}_{ijk}, \hat{s}_{ij}, \hat{w}_{ij}$  are defined as follows

$$\begin{aligned} \hat{p}_{ij} &= (\lambda_j - \lambda_i) r_j \bullet r_i + r_j \bullet A r_i - r_i \bullet A r_j, \\ \hat{q}_{ijk} &= -(r_k \bullet \mu_j) r_j \bullet r_i - \mu_j (r_k \bullet r_j) \bullet r_i + 2(r_k \bullet \mu_i) r_j \bullet r_i + \mu_i (r_k \bullet (r_j \bullet r_i)) \\ &\quad + (r_k \bullet r_i) \bullet B r_j - (r_k \bullet r_j) \bullet B r_i + r_i \bullet B (r_k \bullet r_j) - r_j \bullet B (r_k \bullet r_i) \\ &\quad + (r_j \otimes r_i) : D^2 B r_k - (r_j \otimes r_k) : D^2 B r_i, \\ \hat{s}_{ij} &= 2\mu_i (r_j \bullet r_i) + r_i \bullet B r_j - r_j \bullet B r_i, \\ \hat{w}_{ij} &= (\mu_i - \mu_j) r_j \bullet r_i - r_j \bullet B r_i + r_i \bullet B r_j. \end{aligned}$$

By using interaction estimates (Lemma 3 and 6) we can obtain

$$\|h(t)\|_{L^1} \leq \frac{\|h(\hat{t})\|_{L^1}}{2} \text{ for all } t > \hat{t}. \quad (73)$$

Combining the above inequality with Proposition 2.2, we have

$$\|h(t)\|_{L^1} \leq L_3 \|h_0\|_{L^1} \leq L_3 \|\bar{u} - \bar{v}\|_{L^1} \text{ for all } t > 0. \quad (74)$$

To translate this result to the  $L^1$  stability estimate of two viscosity solutions  $u, v$  we use the homotopy method as in Bianchini-Bressan. Let  $u^\theta$  be the solution corresponding to the initial data  $\theta\bar{u} + (1 - \theta)\bar{v}$ . Then let  $h$  be defined as  $h^\theta := \frac{du^\theta}{d\theta}$ . Then for all  $t > 0$  we have:

$$\|u(t) - v(t)\|_{L^1} \leq \int_0^1 \left\| \frac{du^\theta(t)}{d\theta} \right\|_{L^1} d\theta \leq L_3 \|\bar{u} - \bar{v}\|_{L^1}. \quad (75)$$

## Stability Estimates

As claimed in Theorem 1 we want to prove vanishing viscosity limit as  $\varepsilon \rightarrow 0$  for the following Cauchy problem

$$u_t^\varepsilon + A(u^\varepsilon)u_x^\varepsilon = \varepsilon(B(u^\varepsilon)_x)_x \text{ and } u^\varepsilon(0, x) = \bar{u}(x). \quad (76)$$

Note that  $u^\varepsilon(t, x) = u(t/\varepsilon, x/\varepsilon)$  where  $u$  solves the following problem with fix viscosity but scaled initial data,

$$u_t + A(u)u_x = (B(u)_x)_x \text{ and } u(0, x) = \bar{u}(\varepsilon x). \quad (77)$$

Observe that

$$TV(\bar{u}(\varepsilon \cdot)) = TV(\bar{u}(\cdot)).$$



Therefore, we obtain

$$TV(u^\varepsilon(t)) \leq L_1 TV(\bar{u}), \quad (78)$$

$$\|u^\varepsilon(t) - v^\varepsilon(t)\|_{L^1} = \varepsilon \|u(t) - v(t)\|_{L^1} \leq L_2 \|\bar{u} - \bar{v}\|_{L^1}, \quad (79)$$

$$\|u^\varepsilon(t) - u^\varepsilon(s)\|_{L^1} \leq L_3 \left( |t - s| + \sqrt{\varepsilon} |\sqrt{t} - \sqrt{s}| \right), \quad (80)$$

$$|u^\varepsilon(t, x) - v^\varepsilon(t, x)| \leq \alpha_1 \|\bar{u} - \bar{v}\|_{L^\infty} \left( e^{\frac{c_1}{\varepsilon}(\beta_1 t - (x-a))} + e^{\frac{c_1}{\varepsilon}(\beta_1 t + (x-b))} \right). \quad (81)$$

The convergence of  $u^\varepsilon$  as  $\varepsilon \rightarrow 0$  follows from a standard argument with an application of Helly's theorem and the  $L^1$  continuity (80).

For  $\delta_0 > 0$  and compact set  $K \subset \mathcal{U}$ , we consider

$$\mathcal{D}_0 := \{u : \mathbb{R} \rightarrow \mathbb{R}^n; u(-\infty) \in K \text{ and } TV(u) \leq \delta_0\} \quad (82)$$

By considering a smaller domain  $\mathcal{D} \subset \mathcal{D}_0$  which is positively invariant, we can set  $S : \mathbb{R} \times \mathcal{D} \rightarrow \mathcal{D}$ . From (80) and (81) we conclude the time continuity and continuous dependence on initial data for  $S_t$ .

## Uniqueness of the semi group

We first recall the Riemann solver for equation of hyperbolic systems

$$u_t + A(u)u_x = 0, \text{ with } u(0, x) = \begin{cases} u_l & \text{for } x < 0, \\ u_r & \text{for } x > 0, \end{cases} \quad (83)$$

with  $|u_l - u_r|$  is small enough. We consider the  $i$ -th rarefaction curve  $\sigma \mapsto \mathcal{R}_i(\sigma; u_-)$  starting from  $u_- \in \Omega$  which satisfies

$$\frac{d}{d\sigma} \mathcal{R}_i(\sigma; u_-) = r_i(\mathcal{R}_i(\sigma; u_-)) \text{ with } \mathcal{R}_i(0; u_-) = u_-. \quad (84)$$

By using implicit function theorem and with the help of strict hyperbolicity, there exist  $\bar{\lambda}_1 < \dots < \bar{\lambda}_{n-1}$ ,  $\{\sigma_i\}_{i=1}^n$  and  $\{w_i\}_{i=0}^n$  such that

$$w_0 = u_l, w_n = u_r \text{ and } w_i = \mathcal{R}_i(\sigma_i; w_{i-1}) \text{ for } i = 1, 2, \dots, n. \quad (85)$$

Moreover,  $\lambda_i(\mathcal{R}_i(\theta\sigma_i; w_{i-1})) \in (\bar{\lambda}_{i-1}, \bar{\lambda}_i)$  for  $\theta \in [0, 1]$  and  $1 \leq i \leq n$  where  $\bar{\lambda}_0 := -\infty$  and  $\bar{\lambda}_n := +\infty$ . Let us consider scalar flux  $F_i$  corresponding to  $i$ -characteristics defined as follows

$$F_i(\omega) := \int_0^\omega \lambda_i(\mathcal{R}_i(s; w_{i-1})) ds. \quad (86)$$

Let  $z_i$  be the unique entropy solution the following Cauchy problem for scalar conservation laws,

$$z_{i,t} + F_i(z_i)_x = 0, \quad (87)$$

$$z_i(0, x) = \begin{cases} 0 & \text{if } x < 0, \\ \sigma_i & \text{if } x > 0. \end{cases} \quad (88)$$

Now, we can describe the solution to (83) as follows,

$$u(t, x) = \mathcal{R}_i(z_i(t, x); w_{i-1}) \text{ for } \frac{x}{t} \in [\bar{\lambda}_{i-1}, \bar{\lambda}_i] \text{ for all } i = 1, \dots, n. \quad (89)$$

We first consider the Riemann data where  $u_-, u_+$  both lie on  $i$ -rarefaction curve. Since the rarefaction curves are straight lines we can write

$$\bar{u}(x) = u^* + \bar{z}(x)r_i(u^*) \text{ where } u^* = u(-\infty).$$

Consider the flux  $F_i$  defined as in (86). Then we note that since the solution  $u^\varepsilon$  is satisfying (1), we obtain

$$z_t^\varepsilon + F(z^\varepsilon)_x = \varepsilon(\mu_i(u^\varepsilon)z_x^\varepsilon)_x \text{ where } u^\varepsilon = u^* + z^\varepsilon r_i(u^*) \text{ and } z(0, x) = \bar{z}(x). \quad (90)$$

Due to uniform parabolicity, global solution  $z^\varepsilon$  exists and  $z^\varepsilon$  converges to entropy solution  $z$  of (87).

Hence, the limit  $u$  can be written  $u(t, x) = u^* + z(t, x)r_i(u^*)$ , in other words,  $u^\varepsilon$  converges to a solution of the hyperbolic system,  $u$  defined as in (89).

### Lemma (Finite speed of propagation)

Let  $\bar{u}, \bar{v} \in \mathcal{D}$ . Then there exists  $\beta_1 > 0$  such that the following holds for  $a, b \in \mathbb{R}$ ,

$$\int_a^b |S_t(\bar{u}) - S_t(\bar{v})| dx \leq L_4 \int_{a-\beta_1 t}^{b+\beta_1 t} |\bar{u} - \bar{v}| dx. \quad (91)$$

We consider an initial data which is perturbation of a Riemann data  $\bar{u}_{Rie} = u_- \chi_{(-,0)} + u_+ \chi_{(0,)}$  defined as follows

$$\bar{u}(x) := \begin{cases} u_- & \text{if } x < \delta, \\ w_i & \text{if } i\delta < x < (i+1)\delta, \text{ with } 1 \leq i \leq n-1, \\ u_+ & \text{if } x > n\delta. \end{cases} \quad (92)$$

Due to finite speed of propagation, up to a small time  $t_0$ , the waves do not interact with each other and the limit solution when  $\varepsilon$  goes to 0 can be written as

$$u^\delta(t, x) = \mathcal{R}_i(z_i(t, x - i\delta); w_{i-1}) \text{ for } x \in [i\delta + \hat{\lambda}t, (i+1)\delta - \hat{\lambda}t] \text{ when } t \in [0, t_0]. \quad (93)$$

Now, by sending  $\delta \rightarrow 0$ , we can obtain that the viscosity solution  $u^\varepsilon$  converges to solution of the hyperbolic system for Riemann data. Since a Lipschitz continuous semigroup is determined by the local in time behavior for piecewise constant data, this characterizes the limit function  $u$ . Moreover, it also says that for any subsequence  $\varepsilon_k \rightarrow 0$ ,  $u^{\varepsilon_k}$  converges to the same limit. This completes the proof of Theorem 2.

MERCI POUR VOTRE ATTENTION!!!